



同济大学土木工程防灾国家重点实验室、桥梁工程系

# 高等结构动力学

## ——之第五讲

### 多自由度系统动力问题

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# 1. 动力学方程

## 1.1 节点平衡方程

$$f_{I1} + f_{D1} + f_{S1} = p_1(t)$$

$$f_{I2} + f_{D2} + f_{S2} = p_2(t)$$

$$f_{I3} + f_{D3} + f_{S3} = p_3(t)$$

.....

$$f_{Ii} + f_{Di} + f_{Si} = p_i(t)$$

## 1.2 结构平衡方程

$$\{f_I\} + \{f_D\} + \{f_S\} = \{p(t)\}$$



# 1.动力学方程(续)

## 1.3 节点弹性力

$$f_{Si} = k_{i1}v_1 + k_{i2}v_2 + k_{i3}v_3 + \dots + k_{iN}v_N$$

$$\begin{Bmatrix} f_{S1} \\ f_{S2} \\ \cdot \\ f_{Si} \\ \cdot \end{Bmatrix} = \begin{bmatrix} k_{11} & k_{12} & k_{13} & \dots & k_{1i} & \dots & k_{1N} \\ k_{21} & k_{22} & k_{23} & \dots & k_{2i} & \dots & k_{2N} \\ \dots & \dots & \dots & \dots & \dots & \dots & \dots \\ k_{i1} & k_{i2} & k_{i3} & \dots & k_{ii} & \dots & k_{iN} \\ \dots & \dots & \dots & \dots & \dots & \dots & \dots \end{bmatrix} \begin{Bmatrix} v_1 \\ v_2 \\ \cdot \\ v_i \\ \cdot \end{Bmatrix}$$

## 1.4 节点阻尼力

$$f_{Di} = c_{i1}\dot{v}_1 + c_{i2}\dot{v}_2 + c_{i3}\dot{v}_3 + \dots + c_{iN}\dot{v}_N$$

$$\begin{Bmatrix} f_{D1} \\ f_{D2} \\ \cdot \\ f_{Di} \\ \cdot \end{Bmatrix} = \begin{bmatrix} c_{11} & c_{12} & c_{13} & \dots & c_{1i} & \dots & c_{1N} \\ c_{21} & c_{22} & c_{23} & \dots & c_{2i} & \dots & c_{2N} \\ \dots & \dots & \dots & \dots & \dots & \dots & \dots \\ c_{i1} & c_{i2} & c_{i3} & \dots & c_{ii} & \dots & c_{iN} \\ \dots & \dots & \dots & \dots & \dots & \dots & \dots \end{bmatrix} \begin{Bmatrix} \dot{v}_1 \\ \dot{v}_2 \\ \cdot \\ \dot{v}_i \\ \cdot \end{Bmatrix}$$



# 1.动力学方程(续)

## 1.5 节点惯性力

$$f_{Ii} = m_{i1}\ddot{v}_1 + m_{i2}\ddot{v}_2 + m_{i3}\ddot{v}_3 + \cdots + m_{iN}\ddot{v}_N$$

$$\begin{Bmatrix} f_{I1} \\ f_{I2} \\ \cdot \\ f_{Ii} \\ \cdot \end{Bmatrix} = \begin{bmatrix} m_{11} & m_{12} & m_{13} & \cdots & m_{1i} & \cdots & m_{1N} \\ m_{21} & m_{22} & m_{23} & \cdots & m_{2i} & \cdots & m_{2N} \\ \cdots & \cdots & \cdots & \cdots & \cdots & \cdots & \cdots \\ m_{i1} & c_{i2} & m_{i3} & \cdots & m_{ii} & \cdots & m_{iN} \\ \cdots & \cdots & \cdots & \cdots & \cdots & \cdots & \cdots \end{bmatrix} \begin{Bmatrix} \ddot{v}_1 \\ \ddot{v}_2 \\ \cdot \\ \ddot{v}_i \\ \cdot \end{Bmatrix}$$

结构平衡方程：

$$[M]\{\ddot{v}(t)\} + [C]\{\dot{v}(t)\} + [K]\{v(t)\} = \{p(t)\}$$

$k_{ij}$  — 刚度系数； $c_{ij}$  — 阻尼系数； $m_{ij}$  — 质量系数



## 2.无阻尼自由振动

### 2.1自由振动方程

$$[M]\{\ddot{v}(t)\} + [K]\{v(t)\} = \{0\}$$

令： $\{v(t)\} = \{\hat{v}\}\sin(\omega t + \theta)$

$$\{\ddot{v}(t)\} = -\omega^2 \{\hat{v}\}\sin(\omega t + \theta) = -\omega^2 \{v(t)\}$$

得： $-\omega^2 [M]\{\hat{v}\}\sin(\omega t + \theta) + [K]\{\hat{v}\}\sin(\omega t + \theta) = \{0\}$

### 2.2 振动特征方程

$$([K] - \omega^2 [M])\{\hat{v}\} = \{0\}$$

### 2.3 自振频率方程

$$\| [K] - \omega^2 [M] \| = 0$$



## 2.无阻尼自由振动(续)

### 2.4 自振振型分析

令： $[\tilde{E}^{(n)}] = [K] - \omega^2 [M]$

则： $[\tilde{E}^{(n)}] \{\hat{v}_n\} = \{0\}$  (自动满足)

令：

$$\begin{Bmatrix} \hat{v}_{1n} \\ \hat{v}_{2n} \\ \hat{v}_{3n} \\ \vdots \\ \hat{v}_{Nn} \end{Bmatrix} = \begin{Bmatrix} 1 \\ \hat{v}_{2n} \\ \hat{v}_{3n} \\ \vdots \\ \hat{v}_{Nn} \end{Bmatrix}, \quad \begin{bmatrix} e_{11}^{(n)} & | & e_{12}^{(n)} & e_{13}^{(n)} & \cdots & e_{1N}^{(n)} \\ \hline e_{21}^{(n)} & | & e_{22}^{(n)} & e_{23}^{(n)} & \cdots & e_{2N}^{(n)} \\ e_{31}^{(n)} & | & e_{32}^{(n)} & e_{33}^{(n)} & \cdots & e_{3N}^{(n)} \\ \cdots & | & \cdots & \cdots & \cdots & \cdots \\ e_{N1}^{(n)} & | & e_{N2}^{(n)} & e_{N3}^{(n)} & \cdots & e_{NN}^{(n)} \end{bmatrix} \begin{Bmatrix} 1 \\ \hat{v}_{2n} \\ \hat{v}_{3n} \\ \cdots \\ \hat{v}_{Nn} \end{Bmatrix} = \begin{Bmatrix} 0 \\ \cdots \\ 0 \\ \cdots \\ 0 \end{Bmatrix}$$



## 2.4自振振型分析(续)

振型方程分块形式：

$$\begin{bmatrix} e_{11}^{(n)} & [\tilde{E}_{10}^{(n)}] \\ [\tilde{E}_{01}^{(n)}] & [\tilde{E}_{00}^{(n)}] \end{bmatrix} \begin{Bmatrix} 1 \\ \{\hat{v}_{0n}\} \end{Bmatrix} = \begin{Bmatrix} 0 \\ \{0\} \end{Bmatrix}$$

$$[\tilde{E}_{01}^{(n)}] + [\tilde{E}_{00}^{(n)}] \{\hat{v}_{0n}\} = \{0\}$$

多余方程： $e_{11}^{(n)} + [\tilde{E}_{10}^{(n)}] \{\hat{v}_{0n}\} = 0$

振型公式： $\{\hat{v}_{0n}\} = -[\tilde{E}_{00}^{(n)}]^{-1} [\tilde{E}_{01}^{(n)}]$

注意问题：1.振型矢量首项元素取1将影响计算精度

2.较精确的方法是将最大元素单位化



## 2.5 振型矩阵形式

归一化振型矢量：

$$\{\phi_n\} = \begin{Bmatrix} \phi_{1n} \\ \phi_{2n} \\ \phi_{3n} \\ \dots \\ \phi_{Nn} \end{Bmatrix} \equiv \frac{1}{\hat{v}_{kn}} \begin{Bmatrix} 1 \\ \hat{v}_{2n} \\ \hat{v}_{3n} \\ \dots \\ \hat{v}_{Nn} \end{Bmatrix}$$

振型矩阵：

$$[\Phi] = [\{\phi_1\}, \{\phi_2\}, \{\phi_3\}, \dots, \{\phi_N\}] = \begin{bmatrix} \phi_{11} & \phi_{12} & \phi_{13} & \dots & \phi_{1N} \\ \phi_{21} & \phi_{22} & \phi_{23} & \dots & \phi_{2N} \\ \phi_{31} & \phi_{32} & \phi_{33} & \dots & \phi_{3N} \\ \dots & \dots & \dots & \dots & \dots \\ \phi_{N1} & \phi_{N2} & \phi_{N3} & \dots & \phi_{NN} \end{bmatrix}$$





## 2.6 振型正交性质

基本性质：
$$\{\phi_m\}^T [M] \{\phi_n\} = 0 \quad (m \neq n)$$

$$\{\phi_m\}^T [K] \{\phi_n\} = 0 \quad (m \neq n)$$

附加性质：
$$\{\phi_m\}^T [M] ([M]^{-1} [K])^b \{\phi_n\} = 0 \quad (m \neq n)$$

质量矩阵归一化：

令：
$$\{\hat{\phi}_n\}^T [M] \{\hat{\phi}_n\} = 1$$

得：
$$\{\hat{v}_n\}^T [M] \{\hat{v}_n\} = \hat{M}_n$$

则：
$$\{\hat{\phi}_n\} = \{\hat{v}_n\} \cdot \hat{M}_n^{-1/2}$$

$$[\hat{\Phi}]^T [M] [\hat{\Phi}] = [I]$$



## 3.多自由度系统中的叠加方法

### 3.1 正规坐标(Normal Coordinates)

N维系统  $\left\{ \begin{array}{l} \text{N维位移矢量}\{v\} \\ \text{N维振型函数}\{\phi\} \text{——Fourier展开式} \end{array} \right.$

位移表达式：
$$\{v\} = [\Phi]\{Y\} = \{\phi_1\}Y_1 + \{\phi_2\}Y_2 + \cdots + \{\phi_N\}Y_N = \sum_{n=1}^N \{\phi_n\}Y_n$$

正规坐标：
$$Y_n = \frac{\{\phi_n\}^T [M]\{v\}}{\{\phi_n\}^T [M]\{\phi_n\}}, \quad \dot{Y}_n = \frac{\{\phi_n\}^T [M]\{\dot{v}(t)\}}{\{\phi_n\}^T [M]\{\phi_n\}}$$

正规坐标表示优点：

1.  $\{\phi_n\}$ 的正交性使得求解方便
2. 可以采用部分 $\phi_n$ 足够精度描述N维系统



### 3.2 无阻尼系统（解耦方程）

$$\text{振动方程：} [M]\{\ddot{v}(t)\} + [K]\{v(t)\} = \{p(t)\}$$

$$\text{令：} \{\ddot{v}(t)\} = [\Phi]\{\ddot{Y}\}$$

$$\text{得：} [M][\Phi]\{\ddot{Y}(t)\} + [K][\Phi]\{Y(t)\} = \{p(t)\}$$

等式两边左乘  $\{\phi_n\}^T$  可得：

$$\{\phi_n\}^T [M][\Phi]\{\ddot{Y}(t)\} = \{\phi_n\}^T [M][\phi_n]\{\ddot{Y}(t)\} = M_n \{\ddot{Y}(t)\}$$

$$\{\phi_n\}^T [K][\Phi]\{Y(t)\} = \{\phi_n\}^T [K][\phi_n]\{Y(t)\} = K_n \{Y(t)\}$$

$$\{\phi_n\}^T \{p(t)\} = P_n(t)$$

$$\text{正规坐标方程：} \ddot{Y}_n(t) + \omega_n^2 Y_n(t) = \frac{P_n(t)}{M_n} \quad \left( \omega_n^2 = \frac{K_n}{M_n} \right)$$

$$\text{广义参数：} M_n, K_n, P_n$$



### 3.3 粘滞阻尼系统（解耦方程）

振动方程：
$$[M]\{\ddot{v}(t)\} + [C]\{\dot{v}(t)\} + [K]\{v(t)\} = \{p(t)\}$$

广义参数：
$$C_n = \{\phi_n\}^T [c] \{\phi_n\}$$

$$\{\phi_n\}^T [C] \{\phi_m\} = 0 \quad (m \neq n)$$

$$M_n, K_n, P_n$$

正规坐标方程：
$$\ddot{Y}_n(t) + 2\xi_n \omega_n \dot{Y}_n(t) + \omega_n^2 Y_n(t) = \frac{P_n(t)}{M_n}$$

阻尼比：
$$\xi_n = \frac{C_n}{2\omega M_n}$$



### 3.4 正规坐标响应分析

时域方法——Duhamel积分公式

$$Y_n(t) = \int_0^t P_n(\tau) h_n(t - \tau) d\tau$$

$$h_n(t - \tau) = \frac{1}{M_n \omega_{Dn}} \sin \omega_{Dn}(t - \tau) \exp[-\xi_n \omega_n(t - \tau)]$$

频域方法——Fourier变换公式

$$Y_n(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} H(i\bar{\omega}) P_n(i\bar{\omega}) \exp(i\bar{\omega}t) d\bar{\omega}$$

$$H_n(i\bar{\omega}) = \frac{1}{\omega_n^2 M_n} \left[ \frac{1}{(1 - \beta_n^2) + i(2\xi_n \beta_n)} \right]$$

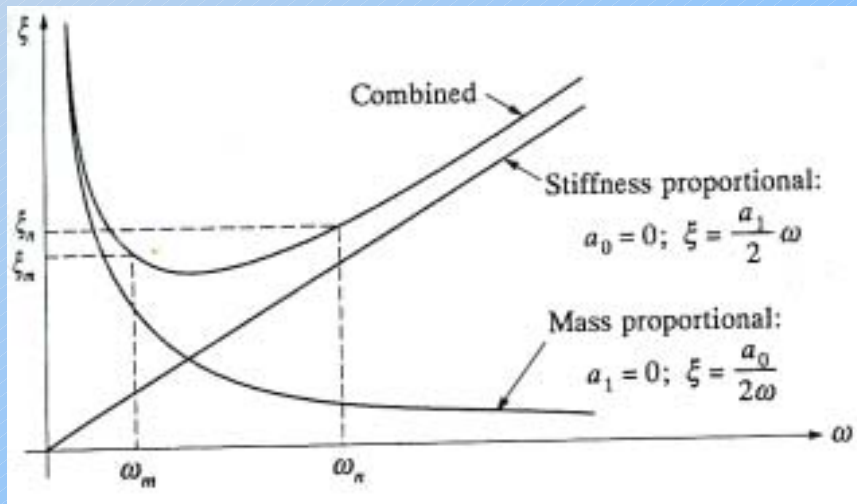
$$P_n(i\bar{\omega}) = \int_{-\infty}^{\infty} P_n(t) \exp(-i\bar{\omega}t) dt$$



### 3.5 线性粘滞阻尼矩阵

瑞利阻尼(Rayleigh Damping) :  $[C] = a_0[M] + a_1[K]$

$$\xi_n = \frac{a_0}{2\omega_n} + \frac{a_1\omega_n}{2}$$



$$\begin{Bmatrix} a_0 \\ a_1 \end{Bmatrix} = 2 \frac{\omega_m \omega_n}{\omega_n^2 - \omega_m^2} \begin{bmatrix} \omega_n & -\omega_m \\ -1/\omega_n & -1/\omega_m \end{bmatrix} \begin{Bmatrix} \xi_m \\ \xi_n \end{Bmatrix}$$

广义瑞利阻尼 :  $[C] = [m] \sum \{a_b\} ([M]^{-1}[K])^b$



## 3.6 耦合振动方程求解法(非粘滞阻尼)

### (1) 时域方法

Dirac delta函数：

$$\delta(t) = \begin{cases} 0 & t \neq 0 \\ \infty & t = 0, \end{cases} \quad \int_{-\infty}^{\infty} \delta(t) dt = 1$$

单个荷载作用：

$$v_{ij}(t) = \int_0^t p_j(\tau) h_{ij}(t - \tau) d\tau$$

$j$ —表示荷载作用

全部荷载作用：

$$v_i(t) = \sum_{j=1}^N \left[ \int_0^t p_j(\tau) h_{ij}(t - \tau) d\tau \right]$$

$i$ —表示结构节点



## (2) 频域方法

动力荷载Fourier变换：

$$P_j(i\bar{\omega}) = \int_{-\infty}^{\infty} P_j(t) \exp(-i\bar{\omega}t) dt$$

单个荷载作用：

$$v_{ij}(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} H_{ij}(i\bar{\omega}) P_j(i\bar{\omega}) \exp(i\bar{\omega}t) d\bar{\omega}$$

全部荷载作用：

$$v_i(t) = \frac{1}{2\pi} \sum_{j=1}^N \int_{-\infty}^{\infty} [H_{ij}(i\bar{\omega}) P_j(i\bar{\omega}) \exp(i\bar{\omega}t) d\bar{\omega}]$$





## 4. 多自由度系统中的分步方法

### 4.1 增量平衡方程

平衡方程：
$$\{\Delta f_I\} + \{\Delta f_D\} + \{\Delta f_S\} = \{\Delta p\}$$

$$\{\Delta f_I\} = \{f_{I_1}\} - \{f_{I_0}\} = [M]\{\Delta \ddot{v}\}$$

$$\{\Delta f_D\} = \{f_{D_1}\} - \{f_{D_0}\} = [C_0]\{\Delta \dot{v}\}$$

$$\{\Delta f_S\} = \{f_{S_1}\} - \{f_{S_0}\} = [K_0]\{\Delta v\}$$

$$\{\Delta p\} = \{p_1\} - \{p_0\}$$

刚度和阻尼系数：
$$k_{ij0} = \left( \frac{df_{S_i}}{dv_j} \right)_0, \quad c_{ij0} = \left( \frac{df_{D_i}}{d\dot{v}_j} \right)_0$$

增量平衡方程：
$$[M]\{\Delta \ddot{v}\} + [C_0]\{\Delta \dot{v}\} + [K_0]\{\Delta v\} = \{\Delta p\}$$



## 4.2 匀加速度方法

伪静力平衡方程： $[\tilde{K}_c]\{\Delta v\} = \{\Delta \tilde{p}_c\}$

$$[\tilde{K}_c] = [K_0] + \frac{2}{h}[C_0] + \frac{4}{h^2}[M]$$

$$\{\Delta \tilde{p}_c\} = \{\Delta p\} + 2[C_0]\{\dot{v}_0\} + [M]\left(\frac{4}{h}\{\dot{v}_0\} + 2\{\ddot{v}_0\}\right)$$

计算步骤： 计算初始 $[\tilde{K}_c]$ 和 $[\Delta \tilde{p}]$

采用高斯分解法求解 $\{\Delta v\}$

计算响应速度 $\{\Delta \dot{v}\} = \frac{2}{h}\{\Delta v\} - 2\{\dot{v}_0\}$

计算初始加速度 $\{\ddot{v}_0\} = [M]^{-1}(\{p_0\} - \{f_{D_0}\} - \{f_{S_0}\})$

注意问题： 线性系统 $[\tilde{K}_c]$ 不变，非线性系统 $[\tilde{K}_c]$ 计算量很大

$[M]^{-1}$ 矩阵必须事先计算并保存好

$\{\Delta \tilde{p}_c\}$ 即使是线性系统每个步长也要重新计算



## 4.3 线加速度方法

伪静力平衡方程：

$$[\tilde{K}_d]\{\Delta v\} = \{\Delta \tilde{p}_d\}$$

$$[\tilde{K}_d] = [K_0] + \frac{3}{h}[C_0] + \frac{6}{h^2}[M]$$

$$\{\Delta \tilde{p}_d\} = \{\Delta p\} + [C_0] \left( 3\{\dot{v}_0\} + \frac{h}{2}\{\ddot{v}_0\} \right) \{\dot{v}_0\} + [M] \left( \frac{6}{h}\{\dot{v}_0\} + 3\{\ddot{v}_0\} \right)$$

计算步骤： 计算初始 $[\tilde{K}_d]$ 和 $[\Delta \tilde{p}]$

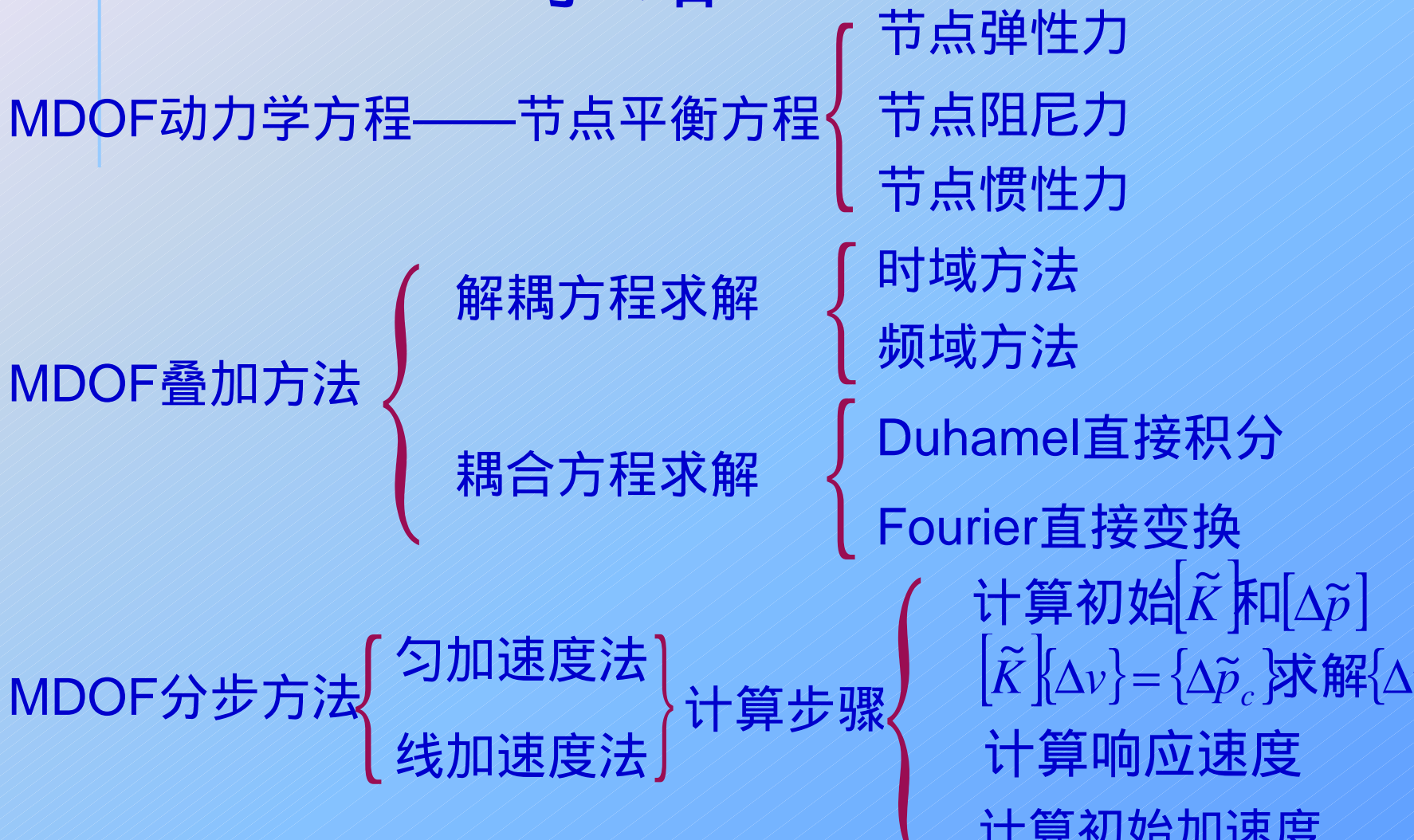
采用高斯分解法求解 $\{\Delta v\}$

$$\text{计算响应速度} \{\Delta \dot{v}\} = \frac{3}{h}\{\Delta v\} - 3\{\dot{v}_0\} - \frac{h}{2}\{\ddot{v}_0\}$$

$$\text{计算初始加速度} \{\ddot{v}_0\} = [M]^{-1} \left( \{p_0\} - \{f_{D_0}\} - \{f_{S_0}\} \right)$$



# 小结





**下周同一时间再见!**